

CSC321: Introduction to Neural Networks and machine Learning

Lecture 16: Hopfield nets and simulated annealing

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Hopfield Nets

- A Hopfield net is composed of binary threshold units with recurrent connections between them. Recurrent networks of non-linear units are generally very hard to analyze. They can behave in many different ways:
 - Settle to a stable state
 - Oscillate
 - Follow chaotic trajectories that cannot be predicted far into the future.
- But Hopfield realized that if the connections are symmetric, there is a global energy function
 - Each “configuration” of the network has an energy.
 - The binary threshold decision rule causes the network to settle to an energy minimum.

The energy function

- The global energy is the sum of many contributions. Each contribution depends on one connection weight and the binary states of **two** neurons:

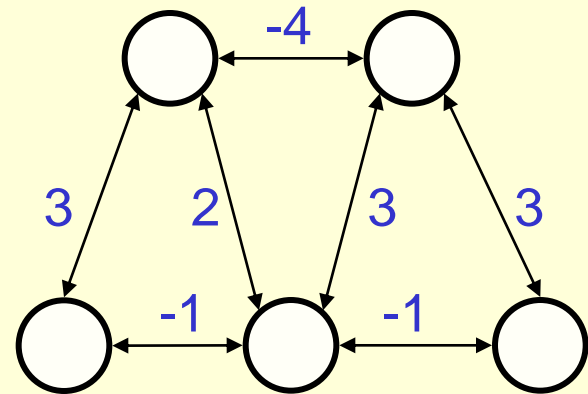
$$E = -\sum_i s_i b_i - \sum_{i < j} s_i s_j w_{ij}$$

- The simple **quadratic** energy function makes it easy to compute how the state of one neuron affects the global energy:

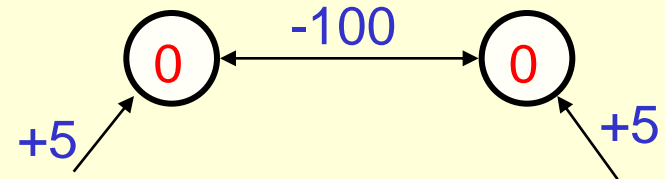
$$E(s_i = 0) - E(s_i = 1) = b_i + \sum_j s_j w_{ij}$$

Settling to an energy minimum

- Pick the units **one at a time** and flip their states if it reduces the global energy.
Find the minima in this net



- If units make **simultaneous** decisions the energy could go up.



How to make use of this type of computation

- Hopfield proposed that memories could be energy minima of a neural net.
- The binary threshold decision rule can then be used to “clean up” incomplete or corrupted memories.
 - This gives a content-addressable memory in which an item can be accessed by just knowing part of its content (like google)
 - It is robust against hardware damage.

Storing memories

- If we use activities of 1 and -1, we can store a state vector by incrementing the weight between any two units by the product of their activities.

$$\Delta w_{ij} = s_i s_j$$

- Treat biases as weights from a permanently on unit
- With states of 0 and 1 the rule is slightly more complicated.

$$\Delta w_{ij} = 4 \left(s_i - \frac{1}{2} \right) \left(s_j - \frac{1}{2} \right)$$

Spurious minima

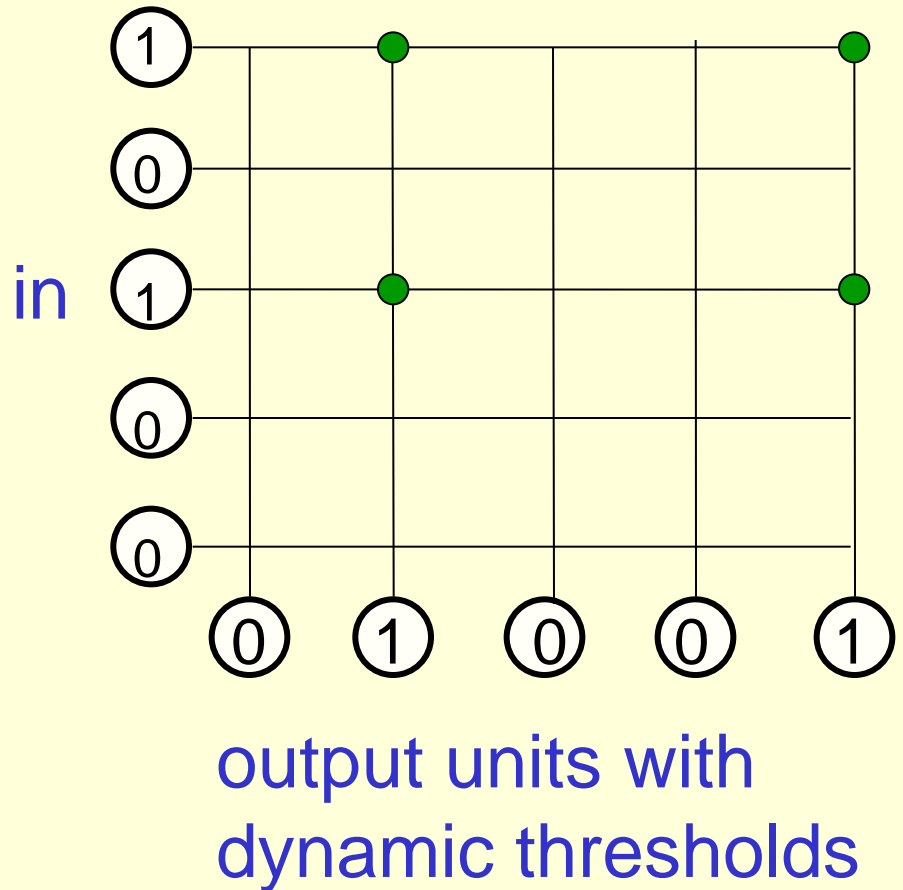
- Each time we memorize a configuration, we hope to create a new energy minimum.
 - But what if two nearby minima merge to create a minimum at an intermediate location?
 - This limits the capacity of a Hopfield net.
- Using Hopfield's storage rule the capacity of a totally connected net with N units is only $0.15N$ memories.
 - This does not make efficient use of the bits required to store the weights in the network.

Avoiding spurious minima by unlearning

- Hopfield, Feinstein and Palmer suggested the following strategy:
 - Let the net settle from a random initial state and then do unlearning.
 - This will get rid of deep , spurious minima and increase memory capacity.
- Crick and Mitchison proposed unlearning as a model of what dreams are for.
 - That's why you don't remember them
(Unless you wake up during the dream)
- But how much unlearning should we do?
 - And can we analyze what unlearning achieves?

Willshaw nets

- We can improve efficiency by using sparse vectors and only allowing one bit per weight.
 - Turn on a synapse when input and output units are both active.
- For retrieval, set the output threshold equal to the number of active input units
 - This makes false positives improbable

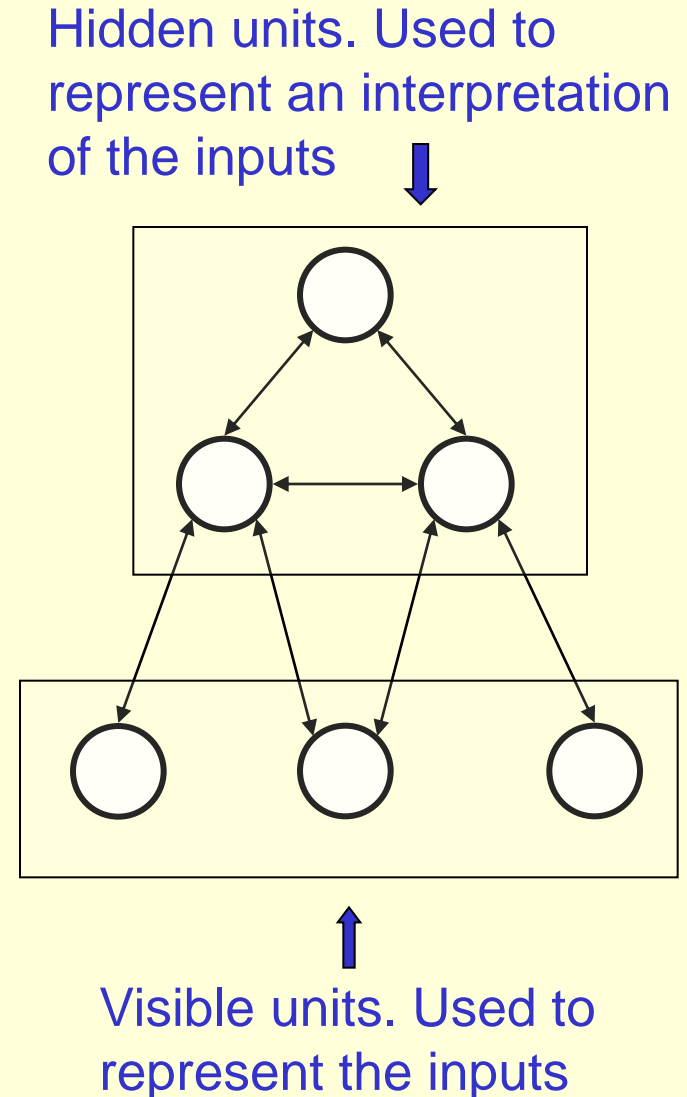


An iterative storage method

- Instead of trying to store vectors in one shot as Hopfield does, cycle through the training set many times.
 - use the perceptron convergence procedure to train each unit to have the correct state given the states of all the other units in that vector.
 - This uses the capacity of the weights efficiently.

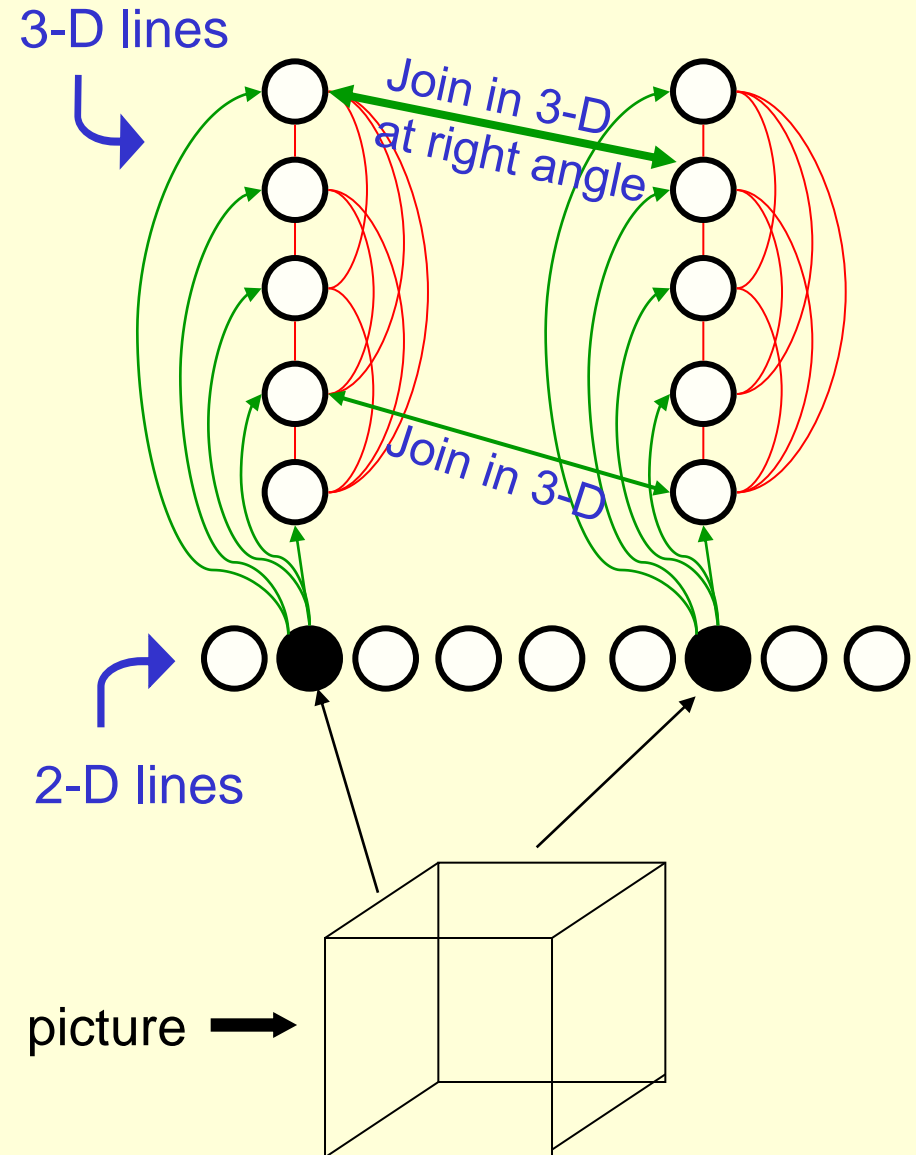
Another computational role for Hopfield nets

- Instead of using the net to store memories, use it to construct interpretations of sensory input.
 - The input is represented by the visible units.
 - The interpretation is represented by the states of the hidden units.
 - The badness of the interpretation is represented by the energy
- This raises two difficult issues:
 - How do we escape from poor local minima to get good interpretations?
 - How do we learn the weights on connections to the hidden units?



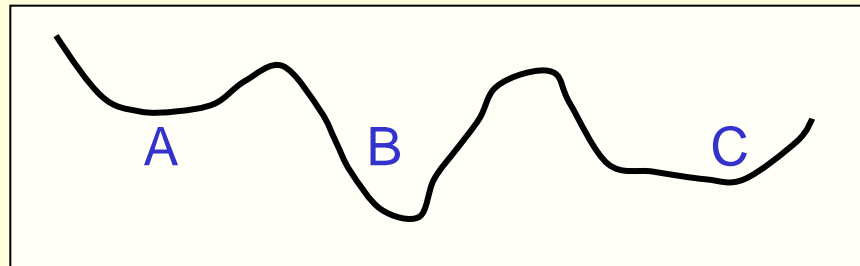
An example: Interpreting a line drawing

- Use one “2-D line” unit for each possible line in the picture.
 - Any particular picture will only activate a very small subset of the line units.
- Use one “3-D line” unit for each possible 3-D line in the scene.
 - Each 2-D line unit could be the projection of many possible 3-D lines. Make these 3-D lines **compete**.
- Make 3-D lines **support** each other if they join in 3-D. Make them **strongly support** each other if they join at right angles.



Noisy networks find better energy minima

- A Hopfield net always makes decisions that reduce the energy.
 - This makes it impossible to escape from local minima.
- We can use random noise to escape from poor minima.
 - Start with a lot of noise so its easy to cross energy barriers.
 - Slowly reduce the noise so that the system ends up in a deep minimum. This is “simulated annealing”.



Stochastic units

- Replace the binary threshold units by binary stochastic units that make biased random decisions.
 - The “temperature” controls the amount of noise
 - Decreasing all the energy gaps between configurations is equivalent to raising the noise level.

$$p(s_i=1) = \frac{1}{1 + e^{-\sum_j s_j w_{ij} / T}} = \frac{1}{1 + e^{-\Delta E_i / T}}$$

↑
temperature

$$\text{Energy gap} = \Delta E_i = E(s_i=0) - E(s_i=1)$$

The annealing trade-off

- At high temperature the transition probabilities for uphill jumps are much greater.

$$p(\textit{pick higher energy state}) = \frac{1}{1 + e^{\Delta E/T}}$$

↑
Energy
increase

- At low temperature the equilibrium probabilities of good states are much better than the equilibrium probabilities of bad ones.

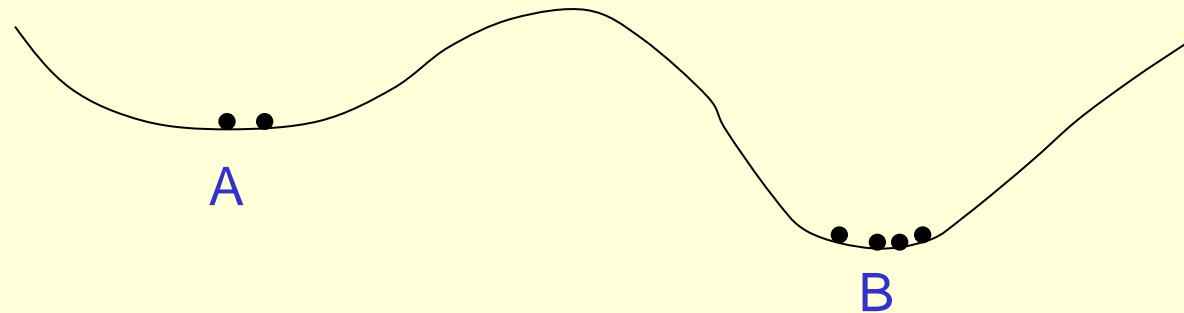
$$\frac{P_A}{P_B} = e^{-\frac{(E_A - E_B)}{T}}$$

How temperature affects transition probabilities

High temperature
transition
probabilities

$$p(A \rightarrow B) = 0.2$$

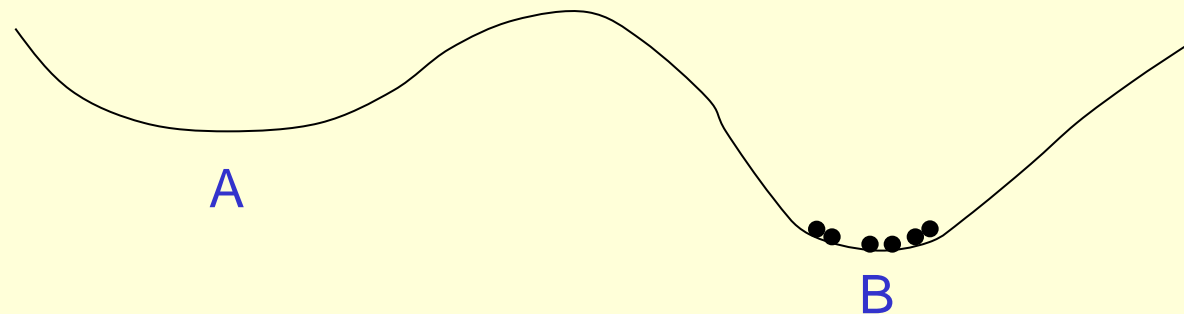
$$p(A \leftarrow B) = 0.1$$



Low temperature
transition
probabilities

$$p(A \rightarrow B) = 0.001$$

$$p(A \leftarrow B) = 0.000001$$



Thermal equilibrium

- Thermal equilibrium is a difficult concept!
 - It does not mean that the system has settled down into the lowest energy configuration.
 - The thing that settles down is the **probability distribution** over configurations.
- The best way to think about it is to imagine a huge ensemble of systems that all have exactly the same energy function.
 - The probability distribution is just the fraction of the systems that are in each possible configuration.
 - We could start with all the systems in the same configuration, or with an equal number of systems in each possible configuration.
 - After running the systems stochastically in the right way, we eventually reach a situation where the number of systems in each configuration remains constant even though any given system keeps moving between configurations

An analogy

- Imagine a casino in Las Vegas that is full of card dealers (we need many more than $52!$ of them).
- We start with all the card packs in standard order and then the dealers all start shuffling their packs.
 - After a few time steps, the king of spades still has a good chance of being next to queen of spades. The packs have not been fully randomized.
 - After prolonged shuffling, the packs will have forgotten where they started. There will be an equal number of packs in each of the $52!$ possible orders.
 - Once equilibrium has been reached, the number of packs that leave a configuration at each time step will be equal to the number that enter the configuration.
- The only thing wrong with this analogy is that all the configurations have equal energy, so they all end up with the same probability.